## 3. Experimental Technique

## a. Apparatus

The electron diffraction tube is a small cathode ray tube (CRT) similar to that in a small (oldfashioned) TV set. The electrons pass through a target that consists of graphite, only a few molecular layers thick, vapor-deposited on a micromesh nickel grid. The target is crystalline in tiny regions, so a "powder" diffraction pattern results, seen as a pair of rings around a central spot on the luminescent screen. See diagram below. The distance between the target and phosphor screen is $L=13.5 \mathrm{~cm}$.


Figure 1 Schematic of the CRT tube

## $\tan (2$ theta) $=\mathbf{r} / \mathrm{D}$ <br> n lambda = 2 d $\sin (t h e t a)$



Figure 7 Crystal Diffraction Geometry
The total angle of a diffraction maxima due to reflection off of a particular crystal lattice plane is twice the angle of incidence on that plane. This angle may be approximated as the distance of the diffraction spot from the undiffracted beam over the distance between the imaging screen and the crystal, i.e. $q \approx 1 / 2 r / D$.

## Multiple diffraction rings


single crystals => diffraction peaks powder => diffraction rings multiple crystals => spotted rings

graphene
also graphite single crystal

fcc aluminum polycrystalline also powder

graphite pyrolytic

## Diffraction pattern for fcc platinum



## Wavelength vs Accelerating Voltage

| $\frac{U_{\mathrm{A}}}{\mathrm{kv}}$ | $\frac{\lambda}{\mathrm{pm}}$ |
| :---: | :---: |
| 4.00 | 19.4 |
| 4.50 | 18.3 |
| 5.00 | 17.3 |
| 5.50 | 16.5 |
| 6.50 | 15.2 |
| 7.00 | 14.7 |
| 7.40 | 14.3 |



Fig. 5: Radii of the first two interference rings as a function of the wavelength of the electrons.

| Direct Lattice | Simple Cubic |  | Body-Centred Cubic | Face-Centred Cubic |
| :---: | :---: | :---: | :---: | :---: |
| Reciprocal Lattice (cube lattice parameter) | Simple Cubic <br> ( $2 \pi / \mathrm{a}$ ) |  | Face-Centred Cubic ( $4 \pi / \mathrm{a}$ ) | Body-Centred Cubic ( $4 \pi / \mathrm{a}$ ) |
| Planes, spacing | (100) | a | No | No |
|  | (110) | $\mathrm{a} / \sqrt{ } 2$ | Yes | No |
|  | (111) | $\mathrm{a} / \sqrt{ } 3$ | No | Yes |
|  | (200) | $\mathrm{a} / 2$ | Yes | Yes |
|  | (210) | $\mathrm{a} / \sqrt{ } 5$ | No | No |
|  | (211) | $\mathrm{a} / \sqrt{ } 6$ | Yes | No |
|  | (220) | $\mathrm{a} / \sqrt{ } 8$ | Yes | Yes |
|  | (221) | a/3 | No | No |
|  | (300) | a/3 | No | No |
|  | (310) | $\mathrm{a} / \sqrt{ } 10$ | Yes | No |
|  | (311) | $\mathrm{a} / \sqrt{ } 11$ | No | Yes |
|  | (222) | $\mathrm{a} / \sqrt{ } 12$ | Yes | Yes |
|  | (320) | $\mathrm{a} / \sqrt{ } 13$ | No | No |
|  | (321) | $\mathrm{a} / \sqrt{ } 14$ | Yes | No |
|  | (400) | $\mathrm{a} / \sqrt{ } 16$ | Yes | Yes |
| Rule: | All h, k, |  | Sum (h+k +1 ) even | $\mathrm{h}, \mathrm{k}, \mathrm{l}$ all odd or all even |
| Ratio of sines of Bragg angles | $\begin{aligned} & 1: \sqrt{2}: \sqrt{ } 3: \\ & \sqrt{8}: \ldots \end{aligned}$ | $: \sqrt{5}: \sqrt{6}$ | $\begin{aligned} & 1: \sqrt{ } 2: \sqrt{ } 3: 2: \sqrt{ } 5: \sqrt{ } 6: \sqrt{ } 7: \\ & \sqrt{8}: \ldots \end{aligned}$ | 1:2/V $3: \sqrt{ } 8 / \sqrt{ } 3: \ldots$ |

## APPENDIX 6

QUADRATIC FORMS OF MILLER INDICES

| Cubic |  |  |  |  | Hexagonal |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $h^{2}+h^{2}+l^{2}$ | Whl |  |  |  |  |  |
|  | Simple | Facecentered | Bodycentered | Diamond | $h^{2}+h h+h^{2}$ | hh |
| 1 | 100 |  |  |  | 1 | 10 |
| 2 | 110 | - • | 110 |  | 2 |  |
| 3 | 111 | 111 | - | 111 | 3 | 11 |
| 4 | 200 | 200 | 200 |  | 4 | 20 |
| 5 | 210 |  |  |  | 5 |  |
| 6 | 211 | - • | 211 |  | 6 |  |
| 7 |  |  |  |  | 7 | 21 |
| 8 | 220 | 220 | 220 | 220 | 8 |  |
| 9 | 300, 221 |  |  |  | 9 | 30 |
| 10 | 310 | -•• | 310 |  | 10 |  |
| 11 | 311 | 311 | - . | 311 | 11 |  |
| 12 | 222 | 222 | 222 |  | 12 | 22 |
| 13 | 320 |  |  |  | 13 | 31 |
| 14 | 321 | -•• | 321 |  | 14 |  |
| 15 |  |  |  |  | 15 |  |
| 16 | 400 | 400 | 400 | 400 | 16 | 40 |
| 17 | 410, 322 |  |  |  | 17 |  |
| 18 | 411,330 | - | 411,330 |  | 18 |  |
| 19 | 331 | 331 | . . . | 331 | 19 | 32 |
| 20 | 420 | 420 | 420 |  | 20 |  |
| 21 | 421 |  |  |  | 21 | 41 |
| 22 | 332 | - - | 332 |  | 22 |  |
| 23 |  |  |  |  | 23 |  |
| 24 | 422 | 422 | 422 | 422 | 24 |  |
| 25 | 500,430 |  |  |  | 25 | 50 |
| 26 | 510,431 | - | 510,431 |  | 26 |  |
| 27 | 511,333 | 511, 333 | . . . | 511, 333 | 27 | 33 |
| 28 |  |  |  |  | 28 | 42 |
| 29 | 520,432 |  |  |  | 29 |  |
| 30 | 521 | - . | 521 |  | 30 |  |
| 31 |  |  |  |  | 31 | 51 |
| 32 | 440 | 440 | 440 | 440 | 32 |  |
| 33 | 522, 441 |  |  |  | 33 |  |
| 34 | 530,433 | - | 530, 433 |  | 34 |  |
| 35 | 531 | 531 | . . . | 531 | 35 |  |
| 36 | 600,442 | 600,442 | 600, 442 |  | 36 | 60 |
| 37 | 610 |  |  |  | 37 | 43 |
| 38 | 611,532 | -•• | 611,532 |  | 38 |  |
| 39 |  |  |  |  | 39 | 52 |
| 40 | 620 | 620 | 620 | 620 | 40 |  |
| 41 | 621, 540, 443 |  |  |  | 41 |  |
| 42 | 541 | . | 541 |  | 42 |  |
| 43 | 533 | 533 | . . | 533 | 43 | 61 |
| 44 | 622 | 622 | 622 |  | 44 |  |
| 45 | 630,542 |  |  |  | 45 |  |
| 46 | 631 | - • | 631 |  | 46 |  |
| 47 |  |  |  |  | 47 |  |
| 48 | 444 | 444 | 444 | 444 | 48 | 44 |
| 49 | 700,632 |  |  |  | 49 | 70, 53 |

(cont.)


Fig. 10-2. Calculated diffraction patterns for various lattices. $s=h^{2}+k^{2}+l^{2}$.
encountered among metals and alloys. The line positions are calculated for $\mathrm{Cu} K \alpha$ radiation, $a=2.50 \mathrm{~A}$, and $c / a=1.633$, which corresponds to the close packing of spheres.)

Powder patterns of cubic substances can usually be distinguished at a glance from those of noncubic substances, since the latter patterns nor-


For our electron and optical diffraction experiments the wavelength is small compared to the lattice constant. Consequently, the wavevector $\mathbf{k}=\mathbf{2} \mathbf{p i} /$ lambda is much longer than the reciprocal lattice spacing $2 \mathbf{p i} / \mathbf{a}$ and the Ewald Sphere cuts through many reciprocal lattice points

reciprocal space
Zero Order Laue Zone
First Order Laue Zone
Second Order Laue Zone
In our experiments, we only see ZOLZ peaks

Ewald sphere geometry for x-ray diffraction


In contrast to electron and optical diffraction where there are many Bragg peaks on the Ewald sphere, in x-ray diffraction there is usually only one Bragg peak on the Ewald sphere

## 2D LATTICES graphene = 1 layer of graphite

e.g. the fused hexagonal pattern of a single layer of GRAPHITE


## Counting Lattice Points/Atoms in 2D Lattices

- Unit cell is Primitive (1 lattice point) but contains TWO atoms in the Motif
- Atoms at the corner of the 2D unit cell contribute only $1 / 4$ to unit cell count
- Atoms at the edge of the 2D unit cell contribute only $\frac{1}{2}$ to unit cell count
- Atoms within the 2D unit cell contribute $\mathbf{1}$ (i.e. uniquely) to that unit cell


$$
d=a \cos (30 \text { degrees })
$$



Figure 9 Lattice Plane Spacing of the
Fundamental Lattice Planes with the Widest Spacing for Hexagonal Crystal Structures

## Analysing a 3D solid

e.g. Graphite $=$ a staggered arrangement of stacked hexagonal layers

Perspective: Clinographic views of solids
Projection onto a Plane: Plan views of solids

## GRAPHITE



## Unit Cell Dimensions

- $\boldsymbol{a}, \boldsymbol{b}$ and $\boldsymbol{c}$ are the unit cell edge lengths


$$
\begin{array}{cc}
a=b \neq c & \text { Atom Positions } \\
\alpha=\boldsymbol{b = 9 0 ^ { \circ }} & (0,0,0) \\
\gamma=120^{\circ} & (0,0,1 / 2) \\
(1 / 3,2,1,1,1 / 3,0)
\end{array}
$$



## Graphite continued



## Unit Cell Dimensions

- $\boldsymbol{a}, \boldsymbol{b}$ and $\boldsymbol{c}$ are the unit cell edge lengths
${ }^{-} \mathbf{a}, \mathbf{b}$ and $\mathbf{g}$ are the angles (a between $b$ and $c$, etc....)


## Counting Atoms in 3D Cells

Atoms in different positions in a cell are shared by differing numbers of unit cells

- Vertex atom shared by $\mathbf{8}$ cells $\mathrm{P} / \mathbf{1}$ atom per cell
- Edge atom shared by $\mathbf{4}$ cells $\mathrm{P} \frac{1}{4}$ atom per cell
- Face atom shared by $\mathbf{2}$ cells $\mathbf{P} / \mathbf{2}$ atom per cell
- Body unique to $\mathbf{1}$ cell $\mathbf{P} 1$ atom per cell



## Aluminum diffraction rings



# DIFFRACTION STANDARD EVAPORATED ALUMINUM <br> PRODUCT NO. 619 



## Aluminum Diffraction Pattern

In order to use a standard in electron diffraction properly, several conditions are necessary:

1) After obtaining the pattern of an unknown, it is necessary to expose the known standard to the same electrical and magnetic conditions, specifically the same lens current or high voltage.
(2) The specimen must be in the same position as the unknown had been. Small movements of the stage are permissible in order to obtain a clearer pattern of the standard.

After developing the diffraction plates, the standard is measured first. The indices are assigned per the attached ASTM "d" spacings. After assigning the spacings, calculate " K " using the formula $\mathrm{K}=\mathrm{Sd}$; where " K " is a constant that represents wave length of the beam, camera length and associated variable crystallographic data, " S " is the diameter of the ring in centimeters and " d " is the interplanar spacing in angstroms. The values for " K " for the first five lines should be within $1 \%$ of each other. Use the mean value of " $K$ ".

Next, the pattern of the unknown is measured, and using "K", determined from the standard, the "d" spacings of the unknown are calculated. It is then necessary to establish identity of the unknown from the ASTM published data.

Aluminum: ASTM "d" spacings

| Miller <br> Indices hk1 | Lattice <br> Spacing $\mathbf{d}(\mathbf{A})$ | Intensity <br> $\mathbf{I}$ | Lattice Constant <br> $\mathbf{a ( \mathbf { A } )}$ |
| :---: | :--- | :---: | :---: |
| 111 | 2.338 | 100 | 4.050 |
| 200 | 2.024 | 47 | 4.048 |
| 220 | 1.431 | 22 | 4.047 |
| 311 | 1.221 | 24 | 4.0489 |
| 222 | 1.1690 | 7 | 4.0495 |
| 400 | 1.0124 | 2 | 4.0496 |
| 331 | 0.9289 | 8 | 4.0490 |
| 420 | .9055 | 8 | 4.0495 |
| 422 | .8266 | 8 | 4.0495 |

Average unit cell from last five lines $\qquad$ .4.0494

## 619 TN 7/02

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